

A High-Finesse Fabry-Pérot Cavity for Hall A

Compton Polarimetry

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Introduction

- Principle of Compton polarimetry
- Compton polarimeter at Jefferson Lab
- Hall A Compton polarimeter upgrade

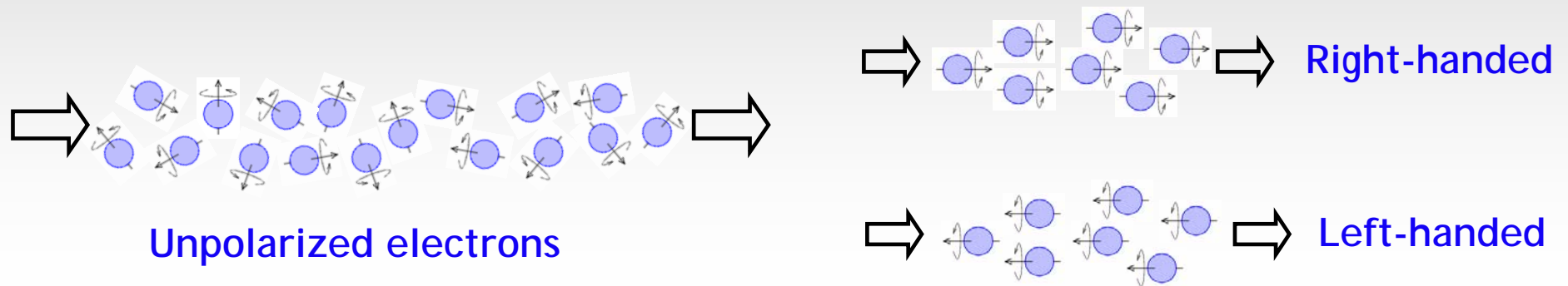
Implementation of Fabry-Pérot cavity

- Cavity optics
- Cavity resonance modes
- Cavity locking

Summary

Polarized Electron Source

- In a **polarized beam**, the spins point in a single direction. The accelerator at Jefferson Lab can produce two types of longitudinally polarized beam

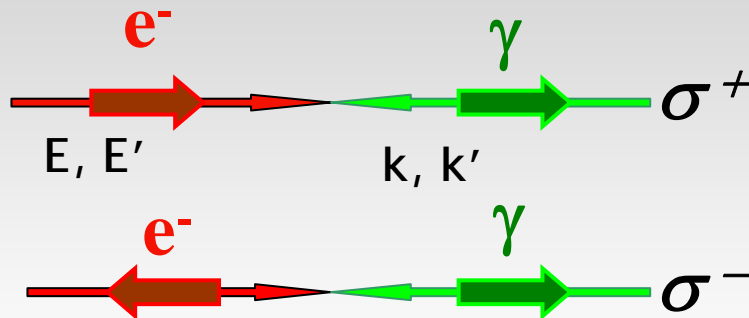


- Via polarized electron scattering, physicists gain insight into what physical processes are involved in an event. This technique is useful for many open research questions, such as:

How densely are neutrons distributed in a heavy nucleus like lead ?
Do strange quarks contribute to the physical properties of protons ?
How is electrical charge distributed inside the neutron ?

- Polarimetry**, the measurement of the beam polarization, allows us to accurately interpret the data and control for errors.

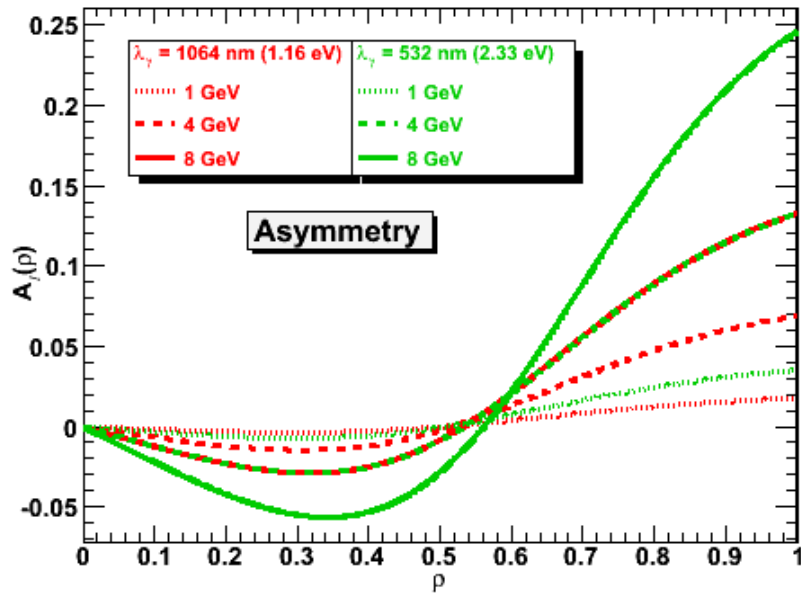
Principle of Compton Polarimetry



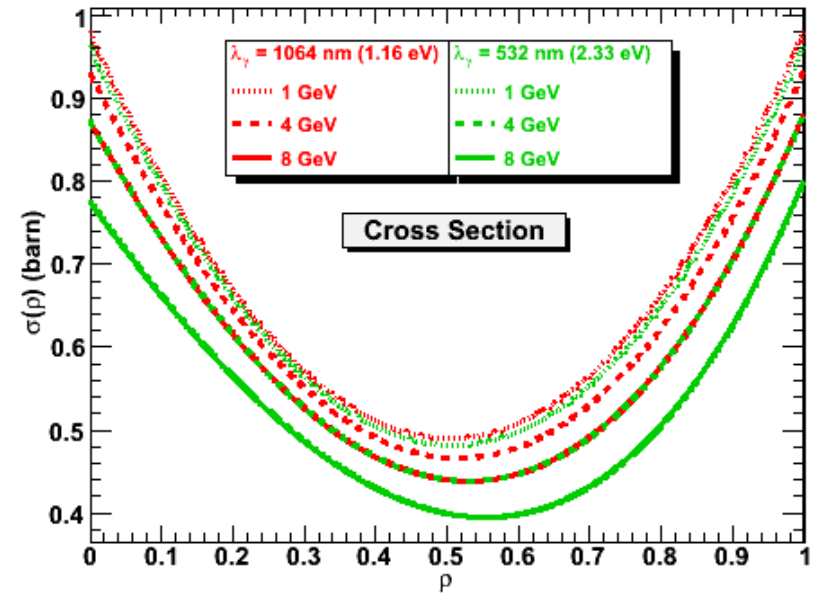
$$A_{\text{exp}} = \frac{n^+ - n^-}{n^+ + n^-} = P_\gamma \times P_e \times \langle A_{th} \rangle$$

Measurable

Theoretical Calculation



$$\frac{\Delta P_e^{stat}}{P_e} \propto \frac{1}{A_l \sqrt{L \sigma \Delta t}}$$



$$\frac{\Delta P_e^{sys}}{P_e} \propto \frac{1}{A_l}$$

Why optical cavity is chosen

$$\sigma_{tot} \approx \frac{8\pi}{3} r_e^2 (1 - x)$$

$$x = \frac{2\gamma\hbar\omega(1 + \beta)}{mc^2} \ll 1$$

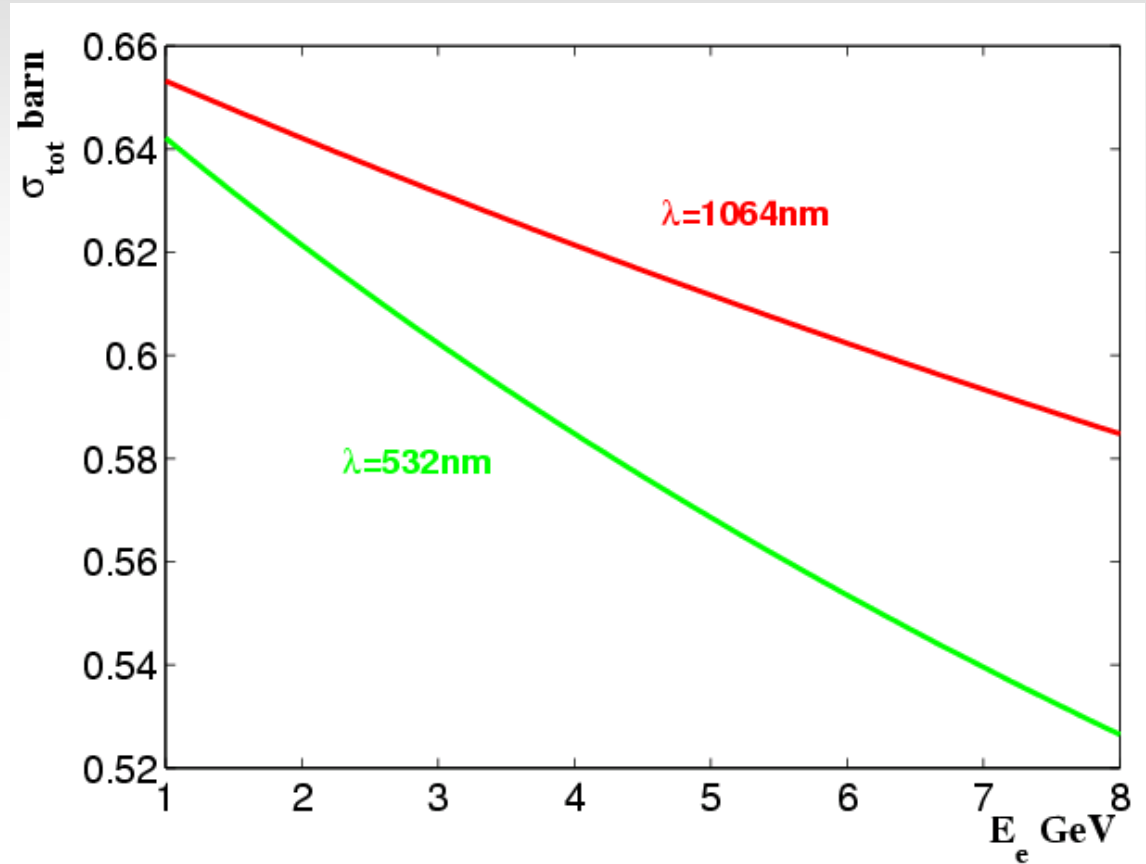
$$r_e = 2.817940325(28) \times 10^{-15} \text{m}$$

$$n = L \times T \times \sigma$$

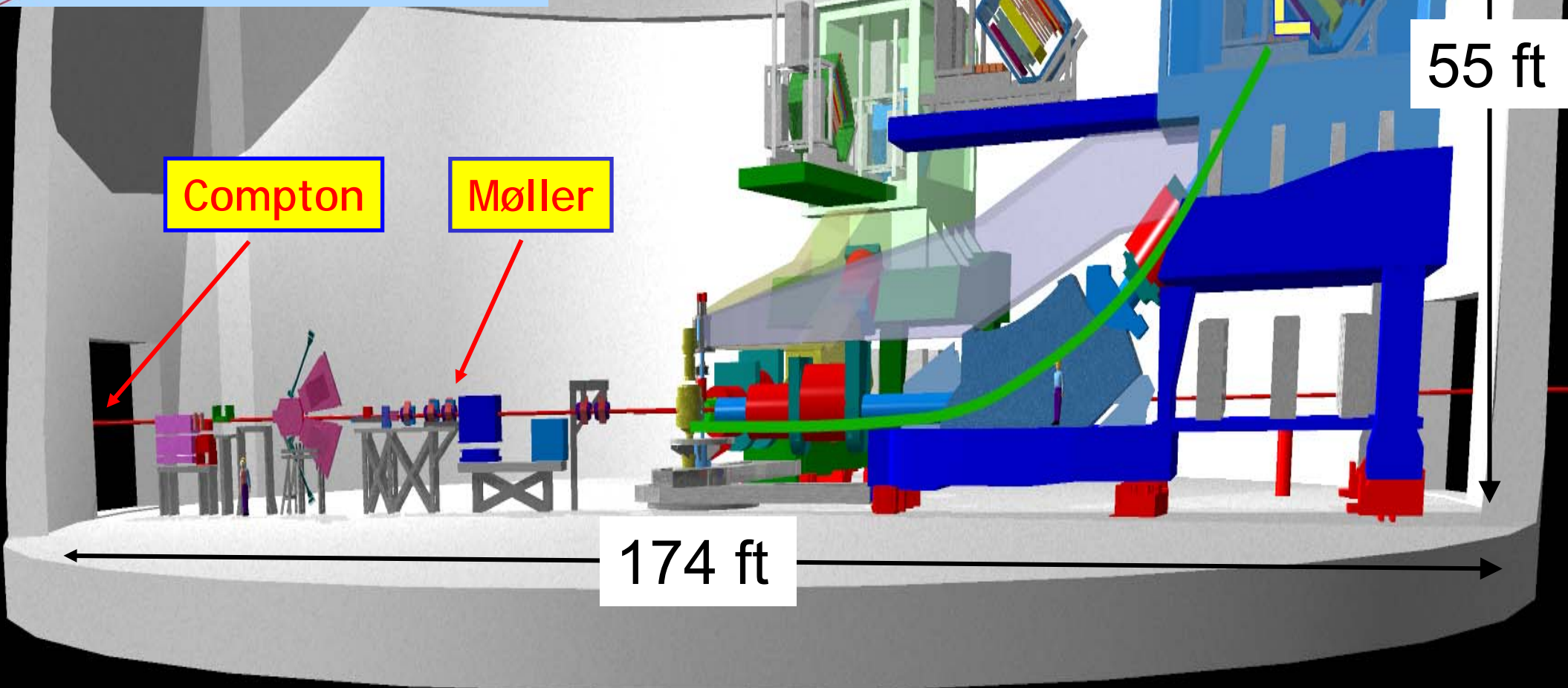
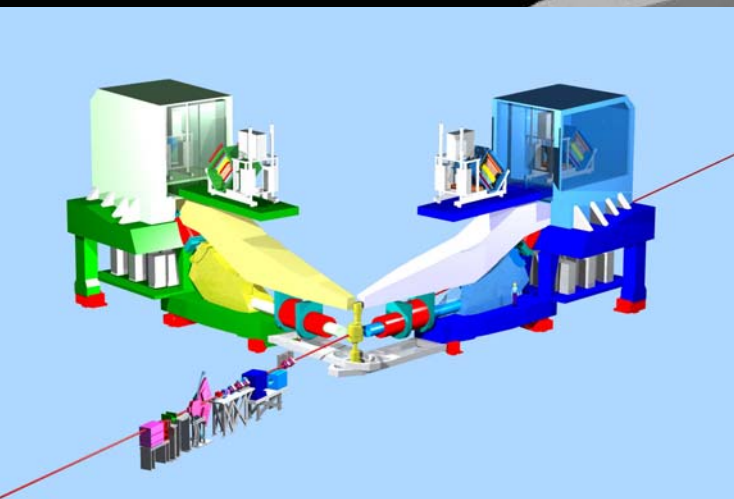
$$L = \iiint v_{rel} \rho_e \rho_\gamma dz dx dy$$

Amplified by optical cavity:

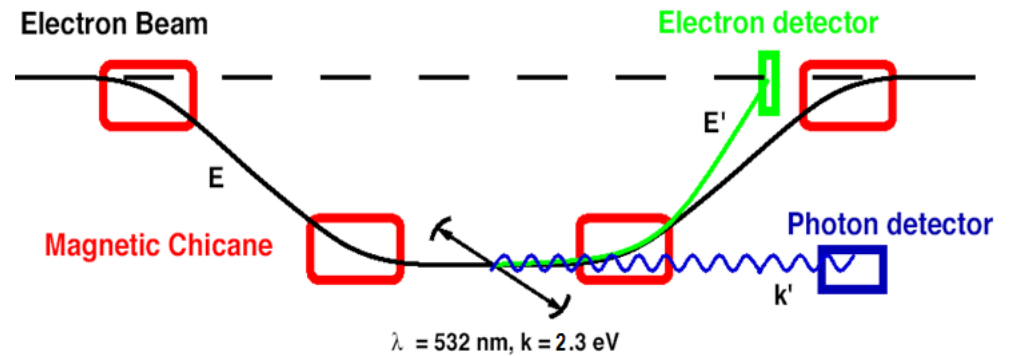
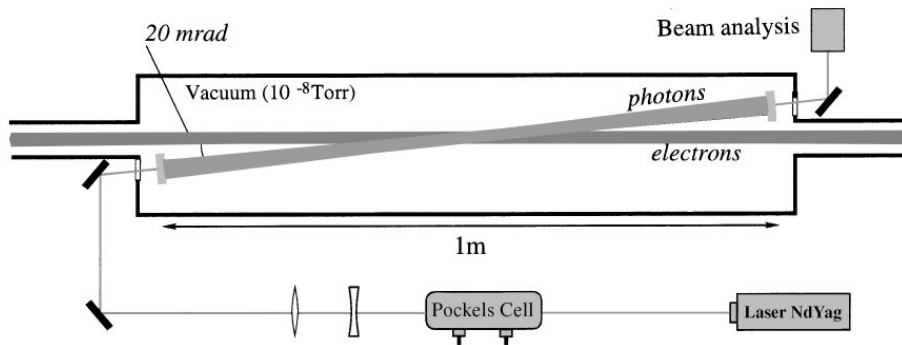
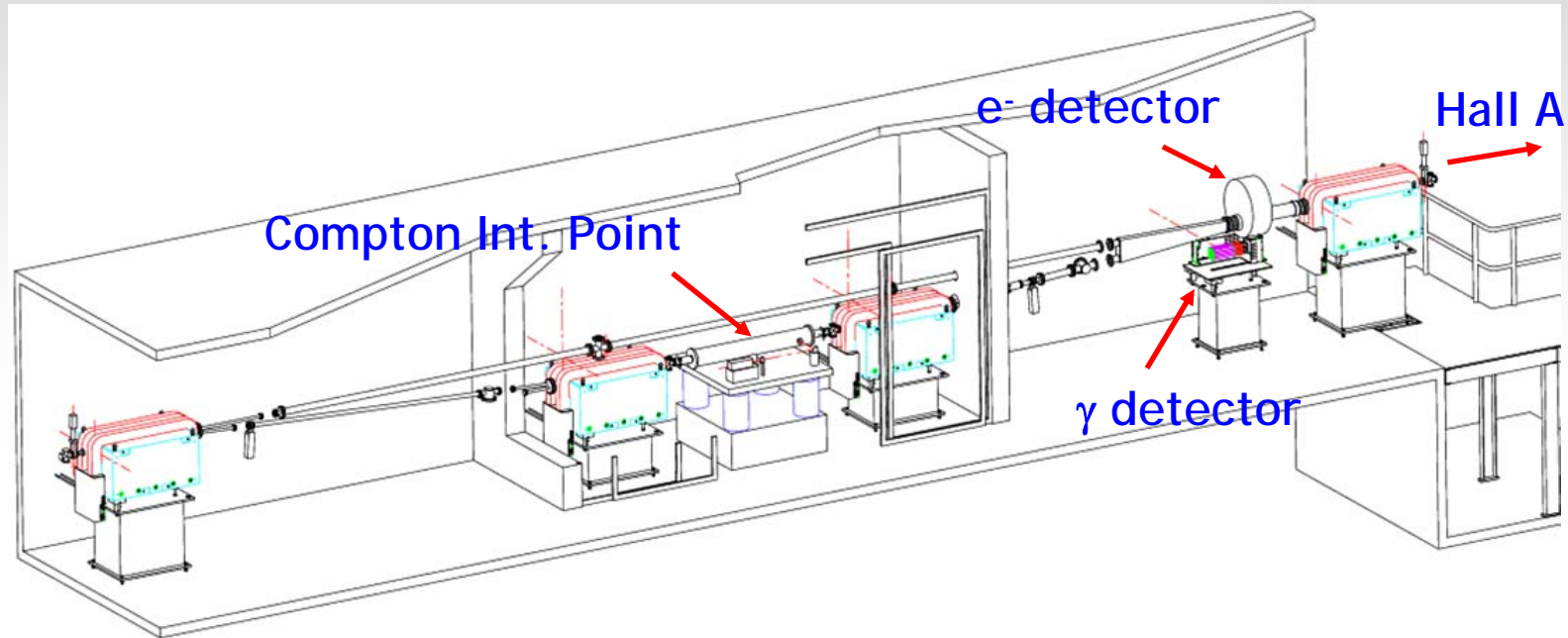
$$L = \mathbf{G} \frac{I_e}{e} \frac{P_L \lambda}{hc^2} \sqrt{\frac{2}{\pi}} \frac{1}{\sqrt{\sigma_e^{y2} + \sigma_\gamma^{y2}}} \frac{1}{\sin \alpha}$$



Compton Polarimeter at Jefferson Lab



Compton Polarimeter at Jefferson Lab



Hall A Compton Polarimeter Upgrade

Motivation:

Improve accuracy of polarization experiments by providing 1% beam polarimetry down to 1 GeV. High precision parity violating experiments (such as PREx) are feasible with this upgrade.

New Electron Detector : High resolution silicon micro strips to improve tracking resolution

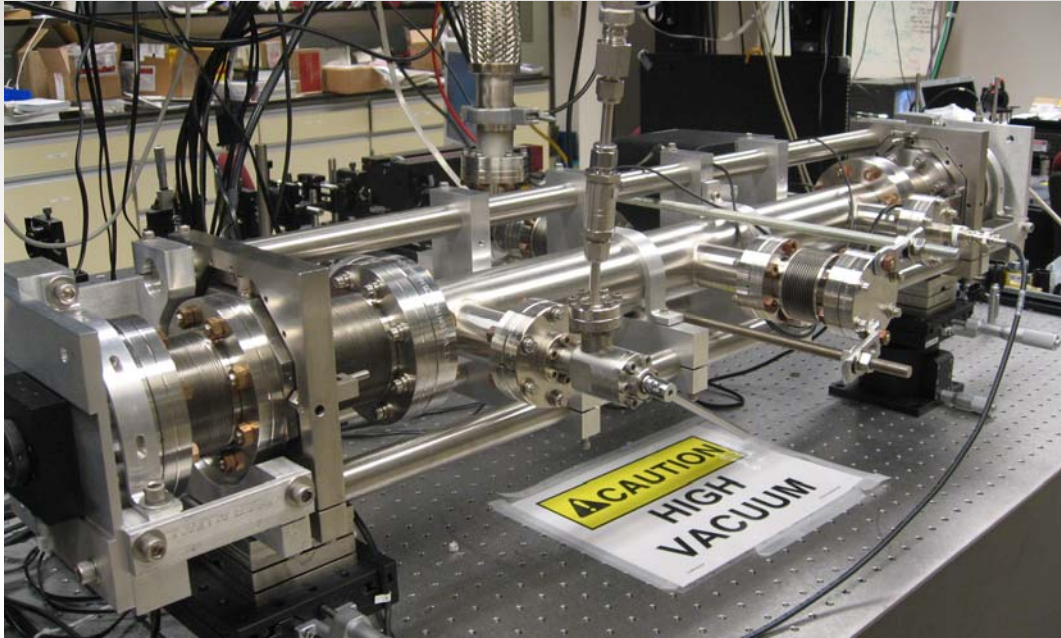
New Photon Detector : Improve systematic uncertainties experienced in the counting method While preserving counting abilities

Green Fabry-Pérot Cavity : Twice the Analyzing power of present IR cavity, two-fold increase in Figure-of-Merit

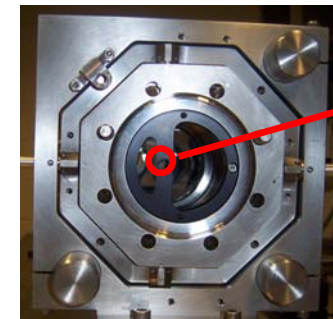
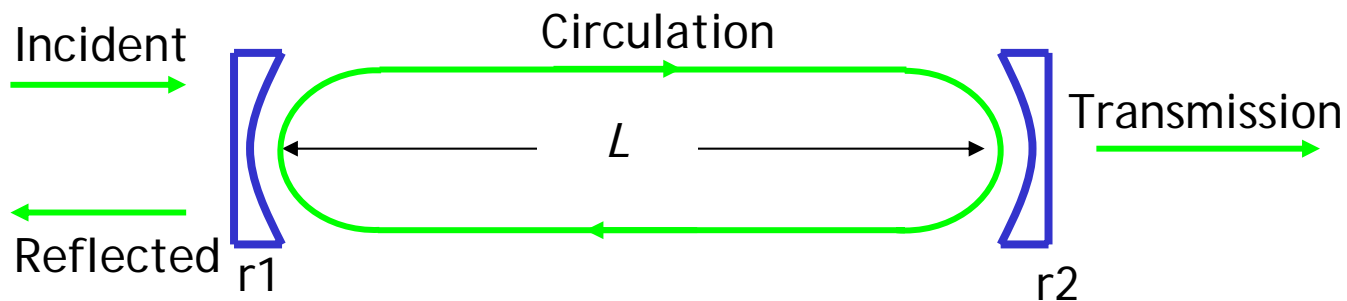
Dynamic Range :

2 GeV ~ 6 GeV → 0.8 GeV ~ 11 GeV

Green Fabry-Pérot Cavity



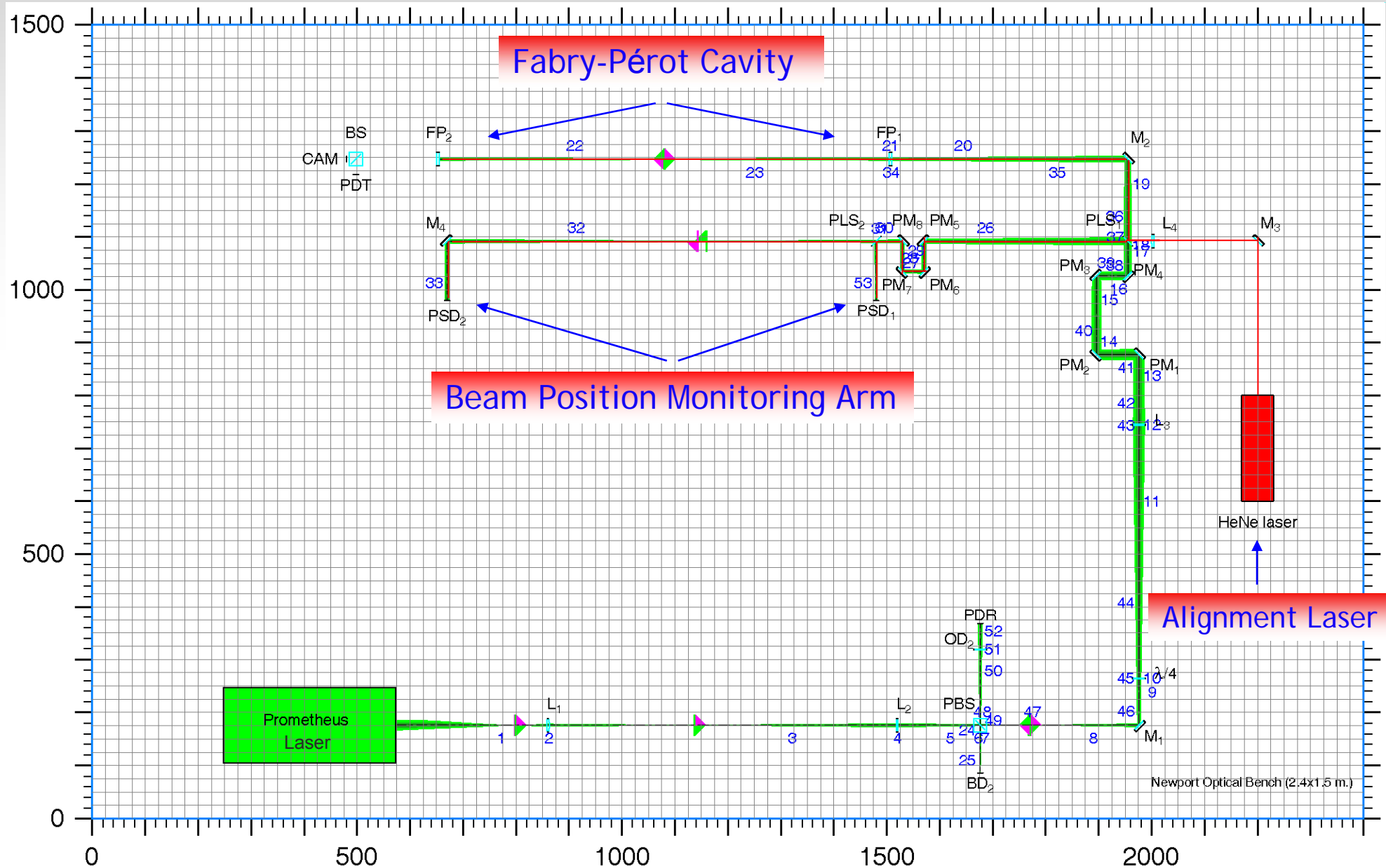
| | |
|----------------------------|-----------------------|
| Wavelength | 532 nm |
| Power | 1,500 Watts |
| Gain | 15,000 |
| Q-factor | 1.8×10^{11} |
| Length | 0.85 m |
| Mode | CW, TEM ₀₀ |
| Free Spectral Range | 176 MHz |
| Cavity Band Width | 3.12 kHz |
| Mirror Reflectivity | 99.996266 % |
| CIP spot size (σ) | 87 μm |



Mounting slot for a cavity mirror

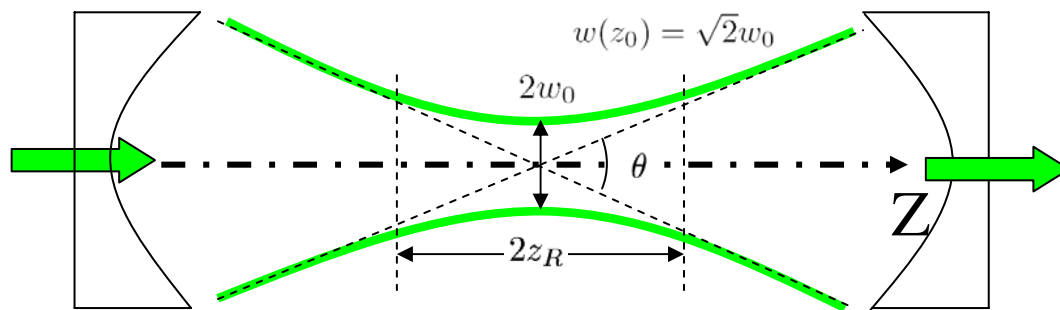
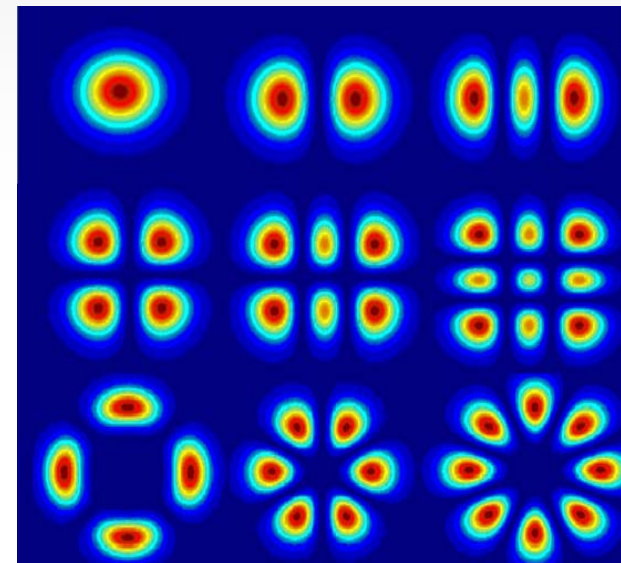
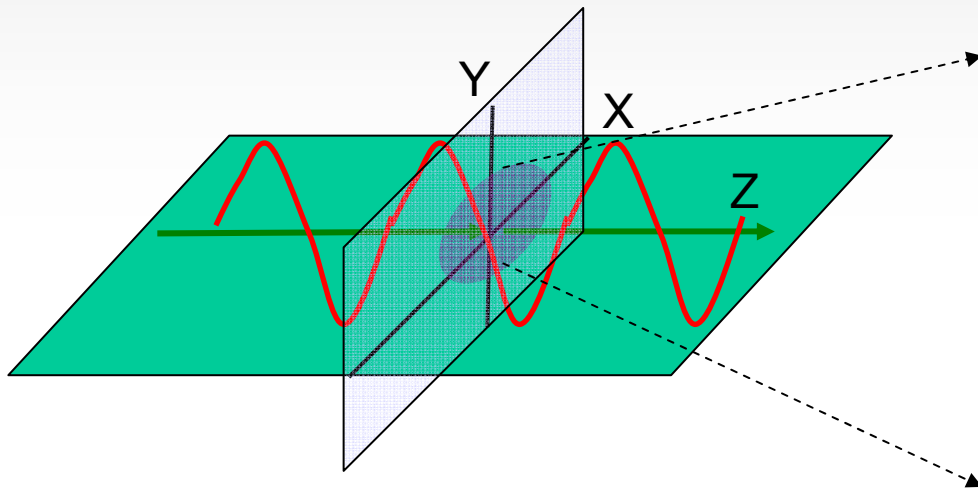
Low power Green Laser -> High Finesse cavity -> Feedback to laser PZT to lock

Optical Setup



Cavity transverse mode

$$E(x, y, z)_{mn} = A \frac{d_0}{d(z)} e^{-ikz} e^{i\omega t} e^{i(m+n+1)\Psi(z)} e^{-ik \frac{x^2+y^2}{2R^2(z)}} e^{-\frac{x^2+y^2}{d^2(z)}} H_m\left(\sqrt{2} \frac{x}{d(z)}\right) H_n\left(\sqrt{2} \frac{y}{d(z)}\right)$$

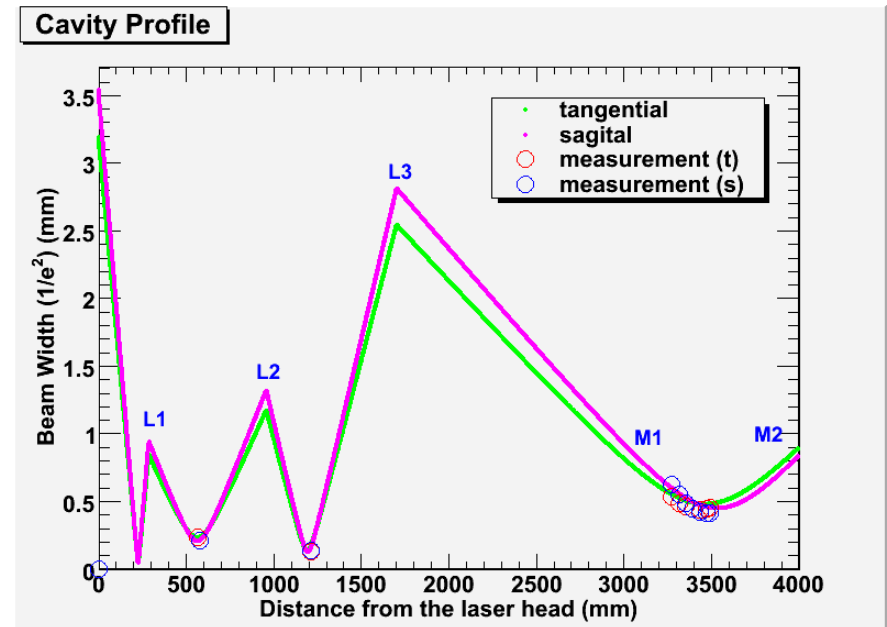
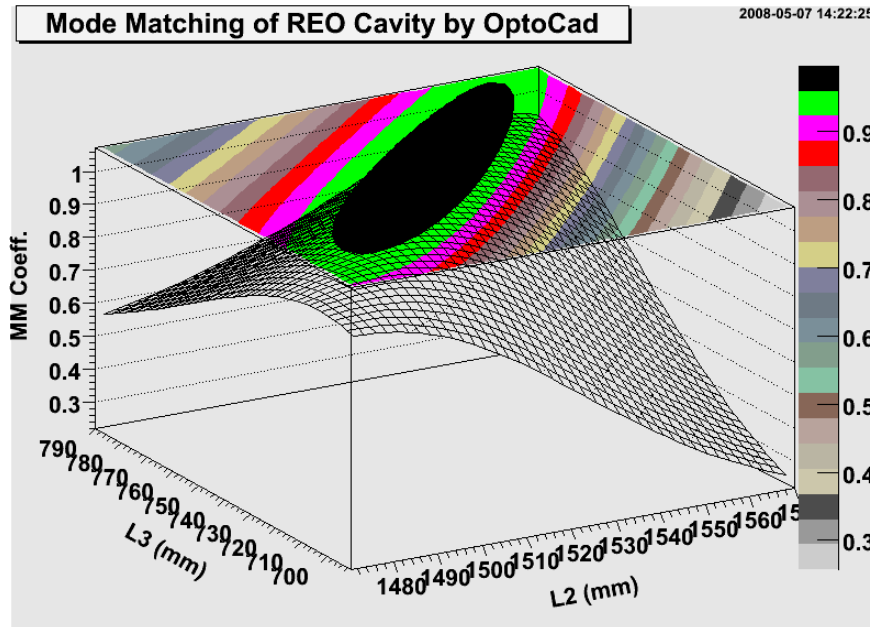
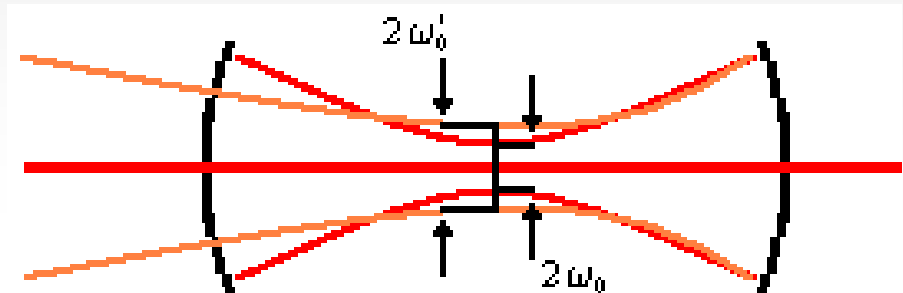
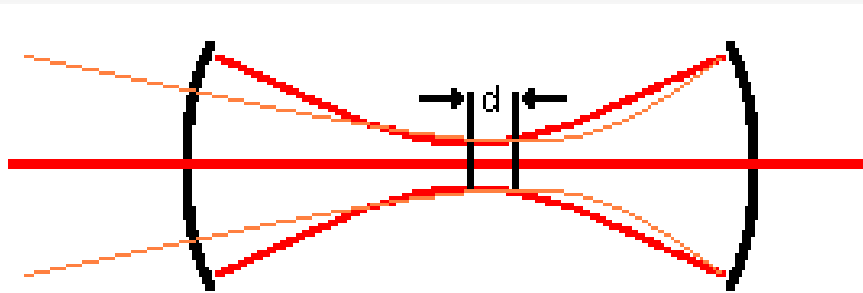


$$w_0^2 = \frac{L\lambda}{\pi} \sqrt{\frac{1+g}{4(1-g)}}$$

$$g = 1 - \frac{L}{R}$$

Cavity Mode Matching

- Laser mode (beam) should match the cavity resonator mode
- Beam waist at the center should match the natural waist of the cavity
- The amount of primary power actually amplified in the fundamental mode

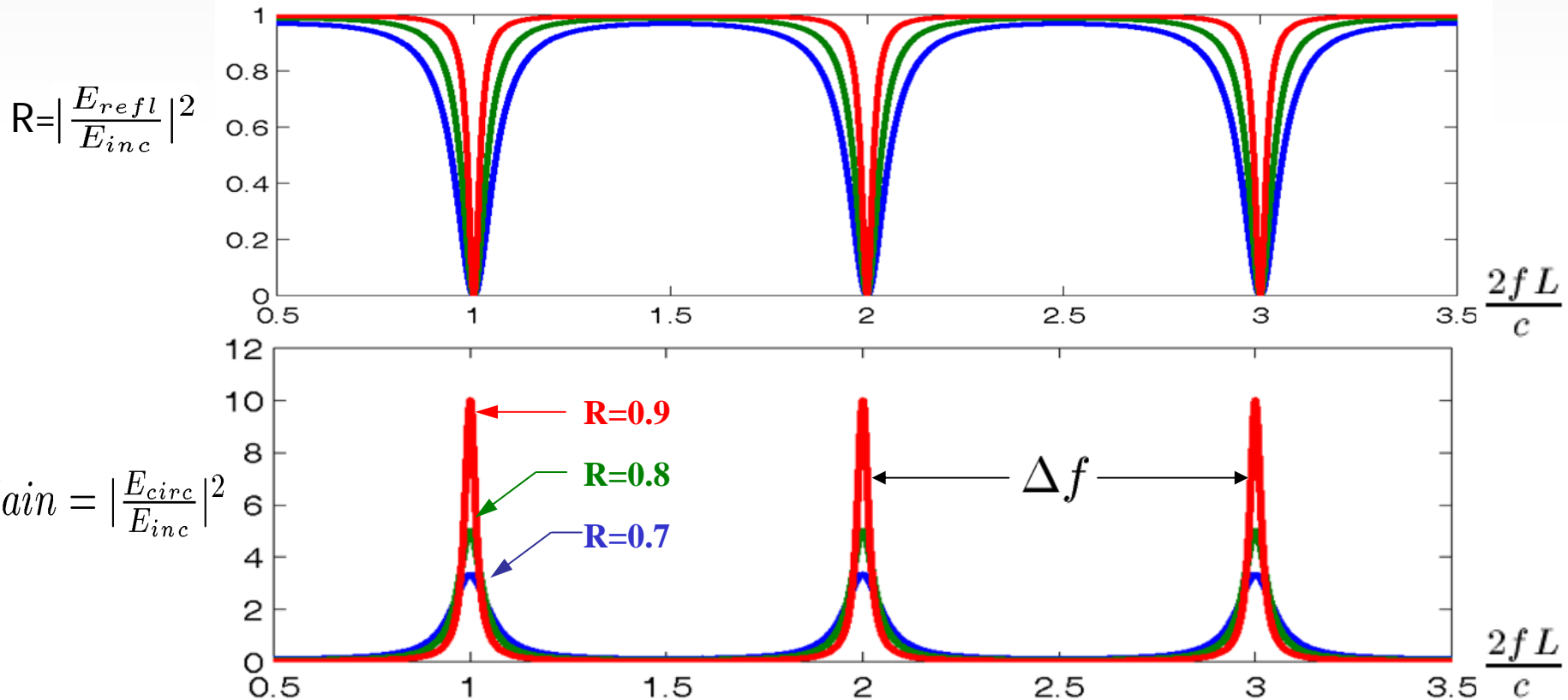
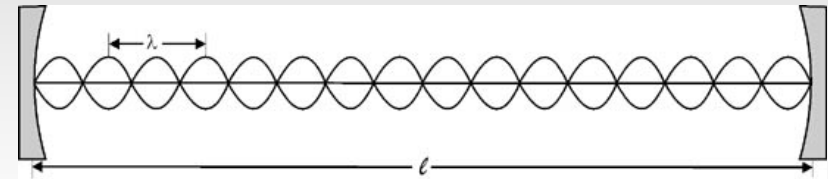


Cavity Resonance

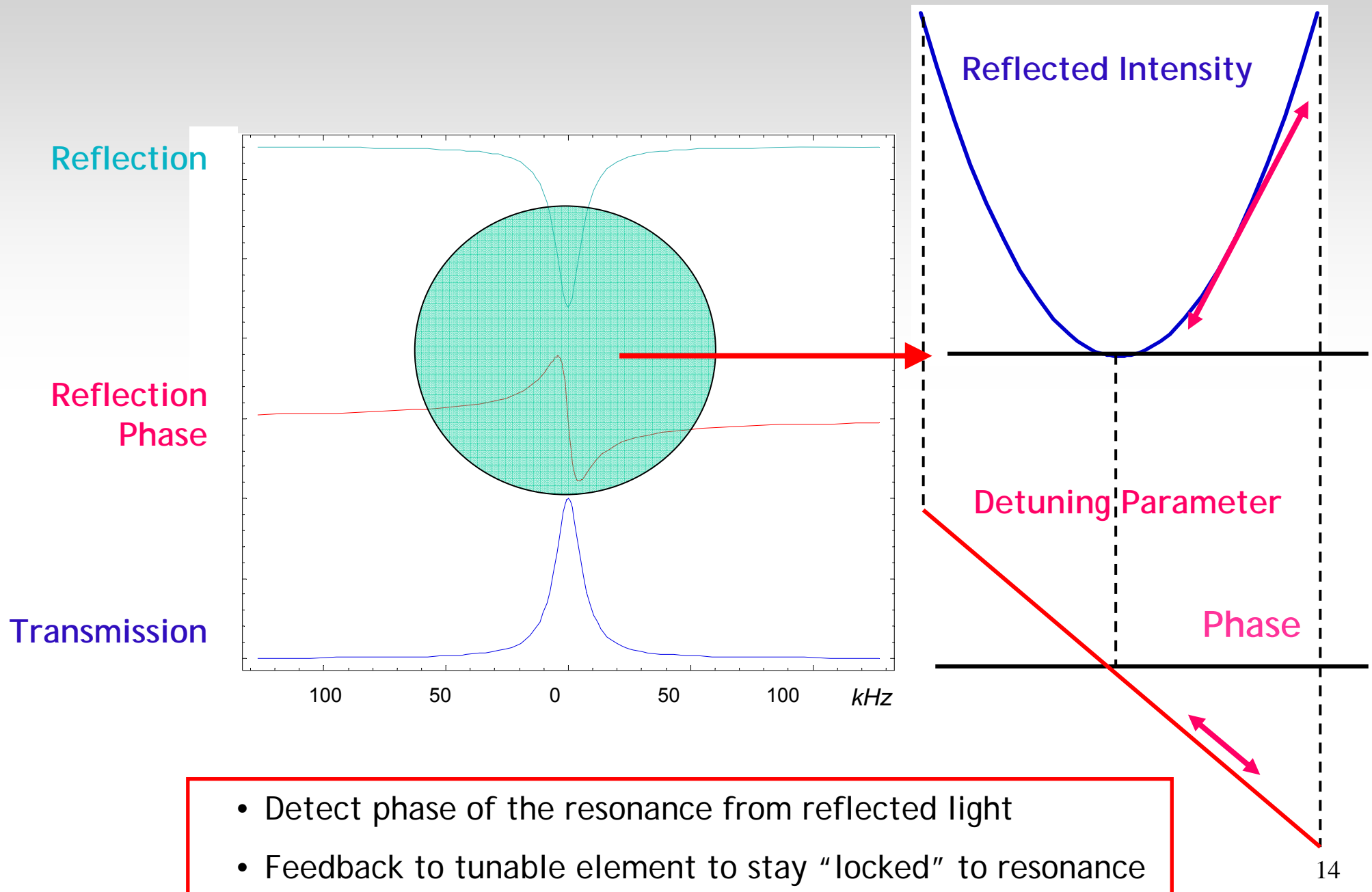
- Keep the cavity resonate forever
- It is very hard to stabilize the cavity length in nm level

$$L = q \frac{\lambda}{2}$$

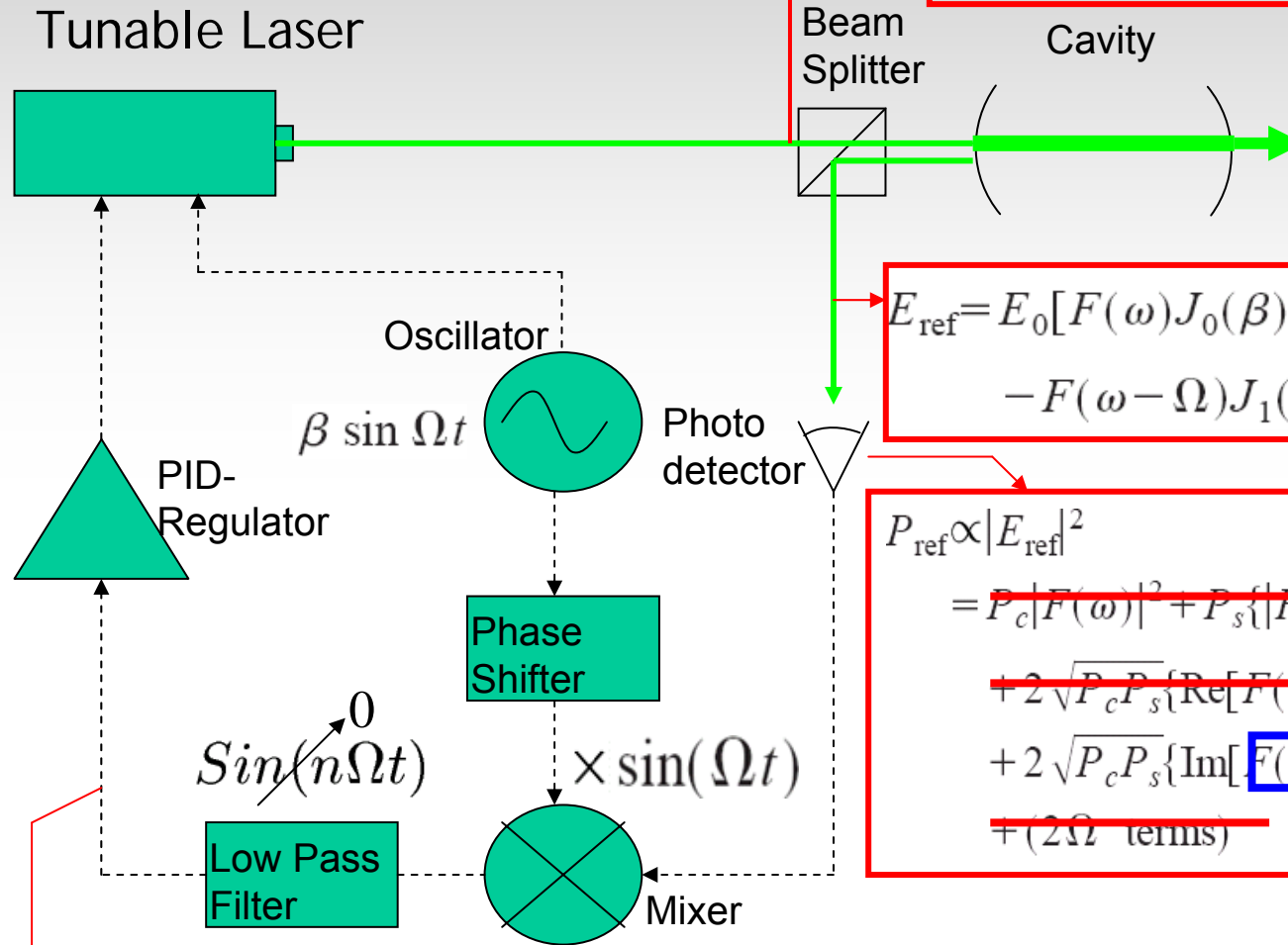
$$\Delta f = \frac{c}{2L}$$



Characteristic Signals of a Fabry-Pérot Cavity



Pound-Drever-Hall Locking Scheme



$$E_{\text{inc}} = E_0 e^{i(\omega t + \beta \sin \Omega t)}$$

$$\approx E_0 [J_0(\beta) + 2iJ_1(\beta) \sin \Omega t] e^{i\omega t}$$

$$= E_0 [J_0(\beta) e^{i\omega t} + J_1(\beta) e^{i(\omega + \Omega)t} - J_1(\beta) e^{i(\omega - \Omega)t}]$$

$$E_{\text{ref}} = E_0 [F(\omega) J_0(\beta) e^{i\omega t} + F(\omega + \Omega) J_1(\beta) e^{i(\omega + \Omega)t} - F(\omega - \Omega) J_1(\beta) e^{i(\omega - \Omega)t}]$$

$$F(\omega) = E_{\text{ref}} / E_{\text{inc}}$$

$$P_{\text{ref}} \propto |E_{\text{ref}}|^2$$

$$= P_c |F(\omega)|^2 + P_s \{|F(\omega + \Omega)|^2 + |F(\omega - \Omega)|^2\}$$

$$+ 2\sqrt{P_c P_s} \{\text{Re}[F(\omega) F^*(\omega + \Omega) - F^*(\omega) F(\omega - \Omega)] \cos \Omega t\}$$

$$+ 2\sqrt{P_c P_s} \{\text{Im}[F(\omega) F^*(\omega + \Omega) - F^*(\omega) F(\omega - \Omega)] \sin \Omega t\}$$

$$+ (2\Omega \text{ terms})$$

$$P_c = J_0^2(\beta) P_0 \quad P_s = J_1^2(\beta) P_0$$

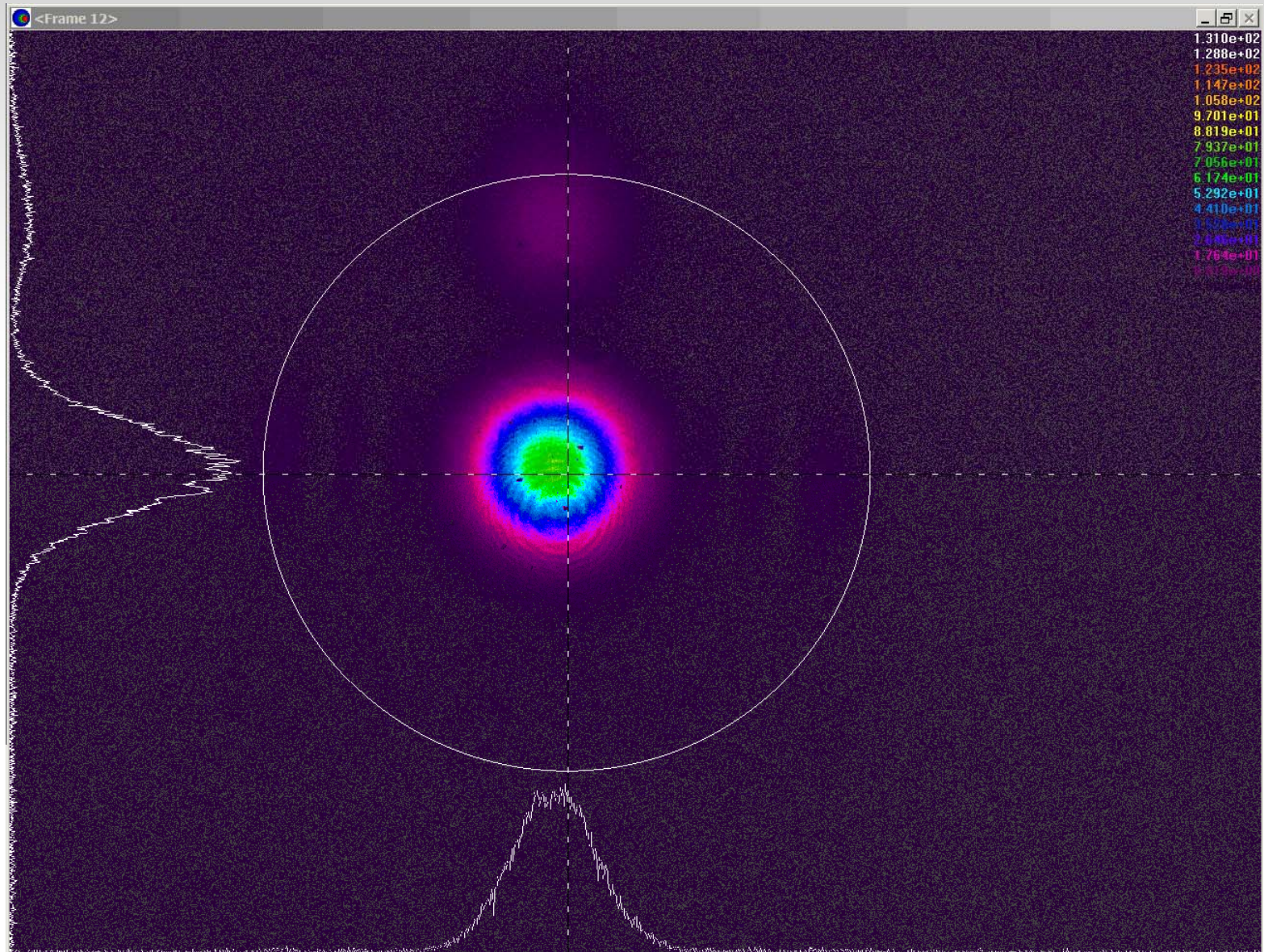
near resonance

$$\epsilon = -2\sqrt{P_c P_s} \text{Im}\{F(\omega) F^*(\omega + \Omega) - F^*(\omega) F(\omega - \Omega)\}$$

$$\Delta\omega_{FSR} / \text{Finesse} \ll \Omega \ll \Delta\omega_{FSR}$$

$$F(\omega) F^*(\omega + \Omega) - F^*(\omega) F(\omega - \Omega) \approx -i2 \text{Im}\{F(\omega)\}$$

Cavity Locking



Summary

- Low gain ($G \sim 1,000$) cavity locking has been accomplished
- Working on design goal cavity ($G \sim 15,000$) locking

Compton Lab ARC L312

